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VIBRATION TESTS OF PRESSURIZED THIN-WALLED CYLINDRICAL SHELLS

by Robert Miserentino and Louis F. Vosteen

Langley Research Center

Langley Station, Hampton, Va.

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SUMMARY

The breathing vibration modes of seven unstiffened cylinders having radius-to-thickness ratios ranging from 324 to 1622 have been determined. Six cylinders were made of stainless steel and one, of aluminum. Their nominal length-to-radius ratio was six. Experimental results are presented to show the variations of resonant frequency and mode shape with internal pressure. The experimental results compared favorably with Reissner's analysis based upon Donnell's equation. However, the theory tended to overestimate the frequency increase which resulted from an increase in internal pressure.

INTRODUCTION

Pressurized thin-walled unstiffened circular cylinders are presently being used as propellant tanks and primary structures for launch vehicles and space-craft. Consequently, their response to the dynamic loads associated with the launch environment is of prime importance to the designer. The solution to dynamic loading problems is dependent on accurate methods for predicting the natural vibration modes and frequencies.

As part of a study to determine the dynamic behavior of thin-walled cylinders, the natural vibration modes and frequencies of seven unstiffened cylinders subjected to internal pressure and various longitudinal loads have been experimentally determined. The radius-to-thickness ratios of these cylinders ranged from 324 to 1624. Six cylinders were made of stainless steel and one, of aluminum. Their nominal length-to-radius ratio was six. The vibration frequencies measured were for the shell breathing modes in which the axis of the cylinder remained undeformed, and the elements of the wall performed harmonic motions predominately in the radial direction. The experimental results were then compared with the results obtained from Donnell's equation (ref. 1), sometimes referred to as Reissner's shallow-shell theory (ref. 2).

SYMBOLS

The units used for the physical quantities in this paper are given both in the U.S. Customary Units and in the International System of Units (SI). Factors relating the two systems are given in reference 3 and those used in the present investigation are presented in the appendix.

```
Α
            ratio of the area that produces a pressure-dependent axial load to the
              total circular cross-sectional area
            mean radius of cylinder, in.
a
            edge fixity constant (c = 0 for simply supported cylinder)
С
            Young's modulus, lb/in<sup>2</sup> (N/m<sup>2</sup>)
E
            frequency, cps (Hz)
f
            skin thickness, in. (m)
h
            frequency parameter, \frac{\rho}{E}(2\pi fa)^2
k
            frequency parameter at zero internal pressure
k_{O}
            free length of cylinder in the longitudinal direction, in.
ı
           number of half-waves in the longitudinal direction
m
           number of full waves in the circumferential direction
n
           nondimensional axial tension due to external load,
n_{x}
           nondimensional axial tension due to internal pressure,
\overline{n}_{\mathbf{y}}
           nondimensional circumferential tension due to internal pressure,
\overline{n}_{m}
           axial load per unit length due to external load, lb/in.
N_x
\overline{N}_{x}
           axial load per unit length due to internal pressure,
             A \frac{pa}{2}, lb/in. (N/m)
\overline{N}_{0}
           circumferential load per unit length due to internal pressure,
              lb/in.
                      (N/m)
```

- p internal pressure, psi (N/m²)
- λ axial wavelength parameter, $\frac{\pi a}{\hbar}(m + c)$
- μ Poisson's ratio
- ρ mass density, $lb-sec^2/in^{l_4}$ (kg/m³)

MODELS, APPARATUS, AND TESTS

Cylinder Models

A total of seven cylinders was constructed; six, of stainless steel and one, of aluminum. Five cylinders had a nominal radius of 6 inches (15 cm) and two had a radius of 4 inches (10 cm). The nominal radius-to-thickness ratios were from 324 to 1624. These ratios are the basis for the reference numbers for each cylinder. Construction details of the cylinders are shown in figure 1, and photographs in figure 2. The photographs, however, are not representative of the condition of the cylinders during vibration tests since they were taken after buckling tests (not reported herein) had been performed. The physical properties and dimensions of each cylinder are given in table I.

Steel end-rings were bonded with an epoxy adhesive to facilitate handling and attachment to the test fixture. In all cases the various end-rings provided essentially clamped supports.

Cylinders 32^{4} , 601, 1001, and 1502 had end-rings that were attached to the support fixture with screws through radial holes. This method caused distortion of the cylinder and made it difficult to obtain a tight seal. The other cylinders were mounted with screws extending longitudinally through holes in the end-rings. (See figs. 1(b) and 1(c).)

Test Apparatus

Fixtures.— All cylinders were mounted on support fixtures that were fundamentally the same. The fixtures consisted of a thick-walled cylindrical support affixed to a heavy mounting base (figs. 3 and 4). The support fits inside the cylinder and the aft end-ring of the cylinder is clamped to an attachment-ring while the forward end of the cylinder is closed with a heavy bulkhead. Some details of this design resulted from aerodynamic considerations which are not reported herein.

For the present vibration tests, the only difference in fixtures or endrings which could affect the results involves the attachment-ring near the base. The cylinders with a 6-inch (15-cm) radius were attached, as shown in figures 3(a) and 3(b), directly to a heavy ring with a T-shaped cross section. The ring moved on three 0 rings that sealed the cylinder chamber. The T-ring provided a means of applying an axial tensile or compressive load to the cylinder. The cylinders with a 4-inch (10-cm) radius (fig. $\mathfrak{Z}(c)$) were attached to a heavy ring that was centered on the central shaft by a thin guiding plate. The ring in turn could be attached to a heavy T-ring through four axial bolts instrumented for measuring axial loads. The chamber was sealed by a flexible rubberized fabric membrane as shown in figure $\mathfrak{Z}(c)$. The method of attachment for the cylinders with the 6-inch (15-cm) radius would be expected to approximate more nearly a fixed-end condition than that of the cylinders with the 4-inch (10-cm) radius.

Cylinder pressurizing system. A pressure-regulating system supplied by a compressed-air source was used to pressurize the cylinders. A Bourdon tube pressure gage was used to measure the internal pressure. The gage and the regulator were tapped into opposite ends of the chamber formed by the cylinder and fixture. A calibration test showed the gage to be accurate within one-half of 1 percent of the scale value, and zero shifts during a test were less than one division of the gage, 0.005 psi (34 N/m^2) .

<u>Vibration equipment.-</u> An electromechanically modulated air-powered loud speaker was used to excite harmonic vibrations of the cylinders. (See fig. 5.) This type of vibration exciter has the advantage of not applying a concentrated force or adding an appreciable mass to the thin walls of the cylinders. The measured acoustic pressure at the test cylinder was 113 dB with a flatness tolerance of 15 dB in the frequency range of 60 to 1000 cps (Hz).

Variations in deflection around the circumference were measured with a variable-reluctance pickup at a point 2/5 the length of the cylinder from the aft end. The variable-reluctance pickup was chosen because it does not contact the model. The pickup was rotated around the cylinder on a motor-driven ring at the rate of 1 rpm. The pickup was automatically positioned radially to the desired distance from the cylinder wall by a servosystem. Thus any variations in circularity or misalinement of the pickup support did not affect the amplitude of the ac signal produced by the vibrating cylinder. Longitudinal variations in the response deflections were measured by either a line of 12 internally mounted reluctance pickups or a hand-held velocity probe with a spring constant of 9 grams force per inch (3.4 N/m).

The ac output of one pickup at a time was displayed along with the shaker oscillator signal on an oscilloscope. The frequency of the pickup signal was read on a frequency meter with an accuracy of 0.5 percent.

Test Procedure

Data were obtained by setting either the chamber pressure or the frequency and varying the other until a peak amplitude response (which was taken to be the resonant point) was seen on the oscilloscope. With both the frequency and pressure held at this resonant point, the pickup was driven around the cylinder to determine the location and number of circumferential half-waves. The mode was identified by a double index - n for the number of waves around the circumference and m for the number of half-waves in the longitudinal direction.

Other observations, such as the relative magnitude of the resonant peak, position of node lines with respect to the welds, spacing of node lines, and phase relations were recorded to help to correlate the data.

In order to minimize the locally induced stresses due to misalinement of the end-rings, the attachments were tightened while the cylinder was vibrating in a low-frequency mode. If the frequency was observed to shift as the bolts were tightened, the attachment was realined until the effect was minimized.

Since the response of a mode is influenced by the orientation of the exciting force field, the loud speaker was moved often to reduce the possibility of modes being overlooked. When it was difficult to pick out a resonance peak, the pickup position, as well as the frequency and the pressure, was varied.

RESULTS AND DISCUSSION

The measured pressures and resonant frequencies for all recorded modes are tabulated in table II for each test cylinder. Also noted in table II are instances when the node lines regularly assumed a particular orientation with respect to the longitudinal seams. In addition, data for a mechanically applied axial load are presented in table III for cylinders 324 and 666.

In table III the loads are given in nondimensional form. It should be pointed out that throughout the text n_X is the nondimensional axial tension load due to external loads only and \bar{n}_X is the nondimensional axial tension load due to internal pressure only, and that $\bar{n}_X = \frac{A}{2} \frac{pa}{Eh}$. Where A for the cylinders with a 6-inch (15-cm) radius was 0.2336 and for the cylinders with a 4-inch (10-cm) radius was 0.3237.

Results for Cylinder With
$$\frac{a}{h} = 324$$

The test results in table II for cylinder 32^{l_4} , which is the one with the smallest value of a/h, are summarized in figure 6. The square of the natural frequency is plotted as a function of internal pressure for modes with one half-wave in the longitudinal direction (m = 1) and for a range of circumferential nodes (n = 2 to 9). Solid straight lines representing a "least-squares" fit through the data points are also shown. The frequency-pressure relation represented by these lines can be expressed in general as $f^2 = c_1 + c_2 p$ which is nondimensionalized to the form

$$k = k_0 + \frac{dk}{d\bar{n}_{\varphi}} \,\bar{n}_{\varphi} \tag{1}$$

where

$$k_0 = 4\pi^2 a^2 \frac{\rho}{E} c_1$$

$$\frac{dk}{d\bar{n}_{0}} = 4\pi^{2}ah\rho c_{2}$$

and

$$\bar{n}_{\varphi} = \frac{pa}{Eh}$$

where c_1 and c_2 are constants determined from the experimental data. From this definition of k (eq. (1)) it follows that k_0 is the nondimensional frequency of the unpressurized cylinder.

Similar conclusions were arrived at theoretically by Reissner in reference 2 where it is shown that values of k_0 and $\frac{dk}{d\vec{n}_0}$ can be written as

$$k_{0} = \frac{\lambda^{1/4}}{(n^{2} + \lambda^{2})^{2}} + \frac{(h/a)^{2}}{12(1 - \mu^{2})}(n^{2} + \lambda^{2})^{2} + n_{x}\lambda^{2}$$
 (2)

and

$$\frac{dk}{d\bar{n}_{\varphi}} = \lambda^2 \frac{d\bar{n}_{x}}{d\bar{n}_{\varphi}} + n^2$$

For the cylinders reported herein, $\frac{dk}{d\overline{n}_{\Phi}}$ can be expressed as

$$\frac{dk}{d\bar{n}_{0}} = \frac{A}{2} \lambda^{2} + n^{2}$$
 (3)

where A accounts for the pressure acting on the attachment-ring, seals, and membrane. The definition of the axial wavelength parameter $\lambda = \frac{\pi a}{l}(m+c)$ is based on reference 4 where it is suggested that the coefficient c be added to m to account for the effective change in length of a simply supported cylinder due to the clamping action on the ends. A value of c=0.3 was found in reference 4 to correlate theoretical results for clamped and simply supported cylinders.

A comparison between the theoretical and experimental frequency is shown in figure 7. The results are plotted in terms of the nondimensional frequency k and the nondimensional pressure parameter \bar{n}_{ϕ} . From the figure it can be seen that, with the exception of n=2, which is a predominantly bending mode, the agreement between theory and experiment is good. Even for the n=2 case the slopes $\frac{dk}{d\bar{n}_{\phi}}$ of the lines are approximately the same; the difference is in the value of the intercept k_{O} .

The difference between theory and experiment is more easily seen in figures 8 and 9. In figure 8 both the theoretical and the experimental values of k_0 are plotted as a function of the mode number n. The experimental values in

figure 8 are the intercepts of the "least-square" lines in figure 7 and represent the frequency of the cylinder with no pressure.

In figure 9, the experimental and theoretical values of $\frac{dk}{d\bar{n}_{\phi}}$ are presented as a function of mode number n. Note that the theoretical contribution to a change in frequency by the axial stress due to pressure (the first term in equation (3)) is negligible for m=1 compared with the contribution of circumferential stress; thus, $\frac{dk}{d\bar{n}_{\phi}}\approx n^2$. From figure 9 it can be seen that except for the case where n=2, the experimentally determined slopes are slightly lower than those calculated theoretically. Differences range from 6 to 15 percent for the modes measured.

Results for Cylinder With Various Values of a/h

The effects of increasing the a/h ratio on the accuracy of theoretically predicted frequency can be seen in figures 10 and 11. In figure 10 the experimental results of k_0 as a function of n are compared with the theoretical results for cylinders with a/h values from 601 to 1624. Although the trends of the experiment and theory agree, a comparison of these results shows that, for modes with n near that of the lowest frequency mode, cylinders with small a/h ratios show a much closer agreement than do cylinders of a/h greater than a 1000. Also note that with the exception of n = 2 and n = 3 the experimental values of k_0 were higher than the calculated ones. Good comparison between theory and experiment for n = 2 or n = 3, however, should not be expected as the Donnell equations, used in the derivation, are applicable only for larger values of n.

Figure 11 indicates that the increase in frequency due to an increase in internal pressure $\frac{dk}{d\bar{n}_\phi}$ is always overestimated by theory. The percentage difference between theoretical and experimental values of $\frac{dk}{d\bar{n}_\phi}$ are fairly uniform for each cylinder. However, there is a slight tendency for the percent difference first to increase and then decrease as n is increased, as is the case for the cylinders tested in reference 5 (a/h = 937 and 3000).

Most of the irregularity in the deviation of the data from the trends shown by the theory can be attributed to an insufficient number of points to average out measuring errors or imperfections in the cylinder. Although no measurements were made of deviations from circularity or variations in skin thickness, visual inspection suggests a ranking (in order of least deviation) of 324, 645, 601, 1624, 666, 1502, and 1001.

From figures 8 to 11, it can be concluded that there is good agreement between theoretical and experimental frequencies for the minimum frequency modes in a range of small values of the pressure parameter \bar{n}_{ϕ} from 0.15 \times 10-5 to 10.00 \times 10⁻⁵. Beyond this range the larger theoretical slope $\frac{dk}{d\bar{n}_{\phi}}$ results in increased separation between theory and experiment. (See eq. (1).)

Results for Modes With m > 1

Within the frequency range of the investigation many circumferential modes with more than one longitudinal half-wave (m > 1) were observed. These modes, however, were often dominated by the nearness of similar modes with one longitudinal half-wave and were therefore difficult to isolate. The greatest number of such modes (with (m > 1)) were observed and recorded for cylinder 645. The results of these observations are shown in figure 12, where ko is plotted as a function of n for the modes with m = 1, 2, and 3. In addition, the corresponding theoretical values, as obtained from equation (2), are shown in the figure. The experimental results agree relatively well with theory for m = 2, but for m = 3 the experimental values of k_0 are considerably lower than those obtained theoretically. Similar results were obtained for all other cylinders tested except for cylinder 1001, which had substantially more imperfections than the other cylinders. In this case, poor agreement between theory and experiment was found even for the m = 2 results. In all cylinders the experimental values of for m > 1 were similar to those for m = 1.

Results for Applied Longitudinal Load

In an effort to obtain the influence of longitudinal load on the frequencies of cylinders, two of the cylinders were subjected to longitudinal loads. The results of these tests are shown in figures 13 and 14. In figure 13, the frequency parameter k_0 is plotted against the mechanically applied axial tension load (n_X) for zero internal pressure for cylinder 32^{l_4} . It can be seen that except for n=2, the axial stresses have no appreciable effect on the frequency. This conclusion is in agreement with the theoretical results shown in equation (2).

Similar results were observed for cylinder 666 (fig. 14) when one resonant mode was followed through elastic compression into the buckling load region. In order to stabilize the shape of the cylinder an internal pressure was maintained.

Experimental Observations

The experimental program involved determining modes and frequencies for different loads at over 1000 data points. Most of the modal shapes investigated had some pressure range where their resonant peak and associated mode shape could be accurately observed. Consequently, in these pressure ranges experimental scatter was small. However, identification of the natural mode was difficult for certain pressure ranges. These difficulties could be attributed to a variety of causes. For example, in certain pressure ranges, several modes occur in a narrow frequency bandwidth and, theoretically, could have the same frequency. In these cases neither mode could be observed accurately.

Another important factor causing experimental scatter is the imperfections existing in the cylinders. The effects of general shell imperfections in the experimentally obtained slope (shown in fig. 11) have already been discussed.

Other types of imperfections such as asymmetry in the cylinder and overlapping of the skin at seam joints also contributed to the experimental uncertainties. In all cases an increase in internal pressure tended to decrease the scattering of data points due to imperfections.

As shown in reference 6, slight asymmetries in a cylinder may resolve a given natural mode of a perfect cylinder into two modes of similar shape but slightly different natural frequencies. This effect has been observed in other investigations (for example, ref. 7) and was encountered in the present tests as shown in figure 10(b) for cylinder 645. In this case the imperfections caused by welded seams caused a regular separation of some modes into two different orientations, one positioned so that a node line runs through each welded seam.

Another factor of some interest is the effect of the bolted attachment of the end-rings of the cylinders. To partially evaluate the importance of this attachment, tests were run on cylinder 324 with the attachment bolts removed from the aft end-ring. The cylinder was mounted vertically. The observed new natural frequencies of the modes are shown in figure 15 along with the theoretical curves obtained for a fixity coefficient c of 0.3 and 0. It can be seen that, by freeing one end of the cylinder, the measured frequencies for the small values of n are considerably reduced and higher values of n may be increased by a slight amount. For modes with n near that of the minimum frequency mode, very little effect was observed.

The exact resonant frequency was sometimes obscured by the fact that the amplitude of vibration did not peak at any distinct frequency, but remained fairly constant over a wide frequency bandwidth. Difficulties also arose when trying to determine a resonant mode shape accurately. In these instances a large amplitude was seen over a substantial region of the cylinder with no apparent node lines; it thus became necessary to count the amplitude peaks. A similar phenomenon was observed in reference 5 and was attributed to waves traveling around the cylindrical circumference. This behavior occurred most often in the cylinders with higher a/h values and sometimes in a wide frequency bandwidth that often included more than one resonant frequency.

A slight nonlinearity in the response of the most responsive modes was seen at times. The peak resonant amplitude was greater and at a higher frequency if the shaker frequency was increased to the resonant frequency than if it were decreased through the same resonant mode.

CONCLUDING REMARKS

The natural modes and frequencies of seven thin-walled circular cylinders were determined experimentally for various values of internal pressure and end loads. The cylinders had a length-to-radius ratio of six and the radius-to-thickness ratios varied from 324 to 1622. A small internal pressure caused a significant increase in resonant frequency over that for zero pressure as well as a reduction in experimental scatter which seemed to be associated with initial imperfections. Changes in end load (tension or compression) caused very little change in frequency over the elastic range covered in this investigation.

Comparison of the experimental results with Donnel shallow-shell theory showed reasonably good agreement except for the mode involving two waves in the circumferential direction. However, the theory tends to overestimate the resonant frequency increase which occurred in this experiment due to an increase in internal pressure.

Langley Research Center,
National Aeronautics and Space Administration,
Langley Station, Hampton, Va., June 4, 1965.

APPENDIX

CONVERSION OF U.S. CUSTOMARY UNITS TO SI UNITS

Factors for converting U.S. Customary Units to the International System of Units (SI) are given in the following table:

Physical quantity	U.S. Customary Unit	Conversion factor (*)	SI unit
Axial load	lb/in.	175.1	newtons/meter (N/m)
Frequency	cps	1	hertz (Hz)
Length	in.	0.0254	meters (m)
Mass density	lb-sec ² /in ⁴	703.07	kilograms/meter ³ (kg/m ³)
Pressure	psi	6.895×10^3	newtons/meter ² (N/m ²)
Radius	in.	0.0254	meters (m)
Thickness	in.	0.0254	meters (m)
Young's modulus	lb/in ²	6.895×10^3	newtons/meter ² (N/m ²)

^{*}Multiply value given in U.S. Customary Unit by conversion factor to obtain equivalent value in SI unit.

Prefixes to indicate multiples of units are as follows:

Prefix	Multiple
giga (G) kilo (k) centi (c) milli (m)	10 ⁹ 10 ³ 10 ⁻² 10 ⁻³

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TABLE I.- PHYSICAL PROPERTIES OF CYLINDERS

a/h	,	a.		ı	1	n.	Material	ρ		E			Number of
.,	in.	m	in.	m	in.	mm	ALCO LAL	lb-s ² /in ⁴	kg/m3	psi	GN/m ²	, μ	seam welds
324	6.01	15.27	36.00	91.44	0.0185	0.4699	17-7 PH stainless steel	0.7149 × 10 ⁻³	7639	29.0 × 10 ⁶	200	0.28	2
601	6.01	15.27	36.00	91.44	.0100	.2540	301 stainless steel	.7408	7916	29.0	200	.32	2
645	4.00	10.16	24.00	60.96	.0062	.1575	301 stainless steel	.7408	7916	29.0	200	.32	2
666	4.00	10.16	24.00	60.96	.0060	.1524	301 stainless steel	.7408	7916	29.0	200	.32	2
1001	6.01	15.27	36.00	91.44	.0060	.1524	2024 aluminum	.2524	2699	10.0	72.3	.32	4
1502	6.01	15.27	36.00	91.44	.0040	.1016	304 stainless steel	.7408	7916	29.0	200	.32	ı
1624	6.01	15.27	38.20	97.03	.0037	.0940	301 stainless steel	.7408	7916	29.0	200	.32	2

TABLE II.- MEASURED MODES, RESONANT FREQUENCIES, AND PRESSURES FOR ALL CYLINDERS

Cylinder	m	n	f		р	Cylinder	m	n	£		р	Cylinder	_	_	f	P	
OJ LIME			cps or Hz	psi	kN/m ²	Cylinder	"	"	cps or Hz	psi	kn/m²	Cylinder	m	n	cps or Hz	psi	kN/m ²
324	1	80 00 00 80 80 80 80 80 80 80 80 80 80 8	387 391 398 400 396	6.00 8.00	0 13.78 41.36 55.15 55.15	601	1	3	252 259 269 280 297	2.00	0 6.89 13.78 27.57 48.26	601	1	10	585 673 744 781	5.00 7.00 9.00 10.00	48.26 62.05
		95 p ⁵	403 402	9.00	62.05 68.94				312 313	9.00	62.05 68.94			11	338	0	٥
		3	245	0	0			4	165 190	0 1.00	0 6.89	645	1	12	453 698	2.00	13.78
			255 262 276 281	4.00 8.00	27.57 55.15 68.94				207 241 255 283	2.00 4.00 5.00	13.78 27.57 34.47	049		3	415 415	1.28 1.72	8.82 11.86
		4	168 193 233	0 2.00 6.00	0 13.78 41.36				283 307 319	9.00	48.26 62.05 68.94		1	83 83	989 415 466	5.66 2.61 7.56	39.03 17.99 52.12
			249 265 160	8.00	55.15 68.94			5	122 169 201		13.78		1 1	81 ₄ 81 ₄	466 415 466	8.18 5.79 8.71	56.40 39.92 60.05
		5	228 256 281 303	4.00 6.00 8.00	27.57 41.36 55.15 68.94				256 276 323 355 373	5.00 7.00 9.00	27.57 34.47 48.26 62.05 68.94		-~- -~- -~-	5 5 a 5 5 a 5	415 587 587 587 587 659	1.19 4.53 5.38 5.49 7.60	8.21 31.23 37.09 37.85 52.40
		6	189 237 274 306 335 355	4.00 6.00 8.00	0 13.78 27.57 41.36 55.15 68.95			6	121 180 226 296 328 377 425	4.00 5.00 7.00	0 6.89 13.78 27.57 34.47 48.26		 1	25 25 5 5 5 25 5 25 25 25 25 25 25 25 25	659 830 830 830 415 466	8.38 15.98 16.22	57.78 110.18 111.84 115.77 32.82 40.75
		7	239 294 335 376 413 429	4.00 6.00 8.00	0 13.78 27.57 41.36 55.15 62.05			7	446 270 349 447 500	2.00 4.00 7.00 9.00	13.78 27.57 48.26 62.05			a6 a6 a6 a6 a6 a6 a6	415 523 587 659 698 830	3.07 5.28 6.84 8.78 10.00 14.31	21.17 36.41 47.16 60.54 68.95 98.67
	-	8	318 370 490 508	8.00 9.00	62.05			8	315 410 450 519	4.00 5.00 7.00	13.78 27.57 34.47 48.26			p6 p6 p6 p6	415 466 523 587 830	3.21 4.22 5.59 7.23 15.01	22.13 29.09 38.55 49.85 103.49
		9		0 2.00 6.00 8.00 10.00	41.36 55.15			9	581 222 303 365 463	9.00 0 1.00 2.00 4.00	0 6.89 13.78 27.57			ь ₇ ъ7	52 3 587 659 698 740	3.81 5.00 6.40 7.34 8.18	26.27 34.47 44.13 50.61 56.40
601	1	2	381 392 394	0 1.00 4.00	0 6.89 27.57				502 592 662	5.00 7.00 9.00	34.47 48.26 62.05					9.38 10.54	64.68 72.67
			395 399 397 398	5.00 7.00 9.00 10.00	48.26 52.05			LO	692 420 540	2.00 4.00	13.78			a ₇	415 523 587 659	2.06 3.62 4.69 6.18	14.20 24.96 32.34 42.61

^aNode lines are on seam welds.

 $^{^{\}rm b}{\rm Node}$ lines are off seam welds.

TABLE II.- MEASURED MODES, RESONANT FREQUENCIES, AND PRESSURES FOR ALL CYLINDERS - Continued

1	İ	ſ	f	:	p _	Ī			f	:	p				f	1	,
Cylinder	m	n	cps or Hz	psi	kN/m ²	Cylinder	m	n	eps or Hz	psi	kN/m²	Cylinder	ш	n	cps or Hz	psi	kN/m ²
645	1	a ₇	698 740 784 830	7.95 9.01	48.61 34.13 62.12 70.33	645	1	12	740 784 988	1.16 1.56 3.66	7.99 10.76 25.24	666	1	ъ2	659 671 669 671	6.20	0 27.92 42.75 45.51
		ъв	329 466 587	0.55 2.03 3.51	3.79 13.99 24.20		2	6	523 587 659 830	3.81 5.49 7.38 12.85	37.85				672 673 675 677	6.90 8.00 8.10 8.30	47.58 55.15 55.85 57.23
			659 698 740 784 830 988	5.16 5.91 6.78 7.73	31.37 35.58 40.75 46.75 52.98 77.22			7	415 659 698 740 784 830	5.60 6.54 7.50 8.61	10.07 38.61 45.09 51.71 59.36 67.98				675 674 670 672 674 679 680	8.90 9.50 10.00 10.20 11.10	59.99 61.36 65.50 68.94 70.33 76.53 85.50
		a.g	466 587 784 830 988	3.32 6.57 7.46	12.13 22.89 45.30 51.44 76.40			8	587 698 740 784 830 988	3.08 4.71 5.51 6.31 7.28	21.24 32.47 37.99 43.51			3	415 435 439 437 447 447	2.00 3.20 3.30 3.40 4.30	13.78 22.06 22.75 23.44 29.65 30.34
	į		415 466 587 659 698 740 784	1.07 2.29 3.17 3.74 4.30	4.62 7.38 15.79 21.86 25.79 29.65			9	587 659 698 740 784 988	1.99 2.91 3.36 4.01 4.57	13.72 20.06 23.17 27.65 31.51 54.33				448 451 458 457 462 463 473	4.50 4.80 5.10 5.40 6.00 6.10	31.03 33.09 35.16 37.23 41.36 42.06 45.16 46.89
		в9	830 988 415 698 740	0.59 3.65 4.17	25.17 28.75			10	587 659 698 740 988	2.40	9.10 13.72 16.55 19.65 41.78				472 472 476 479 486 484	6.90 7.30 7.60 8.00 8.10	47.58 50.33 52.40 55.15 55.85
			784 830 988	5.51	33.16 37.99 56.19		3	6	523 587 659		23.03 34.48 97.16				484 491 490 490	8.30	55.85 57.23 59.86 60.54
		^b 10	415 466 659 784	0.37 .75 2.30 3.69					698 740 932	8.09 9.31 12.35					490 492 494 494	8.80 8.90 9.00	60.68 61.37 62.06 62.06
		^a 10	932 988 466	5.53	38.13 43.71			7	587 659 698 740	5•35 6•24	27.03 36.89 43.02 49.36				496 500 506 513	9.30 9.60	64.12 66.19 70.33
		10	587 659 698	1.45 2.11 2.56	9.99 14.55 17.65				784 830	8.17 9.36	56.33 64.54				513 522	11.28 12.40	77.78 78.60
			7 ⁴ 0 784 988	2.98 3.41	20.55 23.51 42.68			8	740 784 830	6.12	36.89 42.19 49.09			14	367 373 383 394	4.30 4.30	28.96 29.65 29.65
		11	587 698 784	2.42	10.89			9	659 698		18.62 22.48				441 446 464	7.65	33.03 46.89 52.75 54.47
awaa	10 74		988	4.77	32.89			10	698	2.31	15.93				485 525	8.10	55.85 76.53

^aNode lines are on seam welds. ^bNode lines are off seam welds.

TABLE II.- MEASURED MODES, RESONANT FREQUENCIES, AND PRESSURES FOR ALL CYLINDERS - Continued

Cylinder	m	n	f	р	Cylinder	m	l n	f		p	Cylinder	m	n	f		P
OJ III dei			cps or Hz	psi kN/m ²	Jojiinici	"	"	eps or Hz	psi	kn/m ²	OJ ZIMET		"	cps or Hz	psi	kN/m ²
666	1	5	402 424	4.30 29.65 4.70 32.41	666	1	6	703 714	10.90	75.16 75.16	666	1	9	1045	9.00	62.01
			418 466 492 519 507 513 529 529 527 536 548 554 559	4.75 32.75 5.75 39.65 6.80 451.02 7.40 51.02 7.40 52.40 7.60 52.40 7.60 52.40 7.50 55.16 8.10 55.85 8.10 55.85 8.10 55.85 8.70 59.99 8.70 59.99 9.00 62.55			7	579 597 648 656 659 677 673 690 697 701 709 713 736 761 766	4.40 5.30 5.40 5.70 5.60 5.79 6.45 6.60 6.70 6.80 7.40 8.00	27.57 36.54 37.23 38.61 39.30 41.36 43.78 44.47 46.20 46.89 51.02 55.15			10	405 543 569 591 671 731 741 800 835 890 910 910 934 978 1029 1179 1188	2.90 2.90 3.35 4.00 4.75 5.20 5.50 5.80 6.70 9.00	9.31
		6	552 578 567 590 651 651 402 412 415	9.10 62.75 9.50 65.50 9.70 66.88 10.10 69.64 12.40 85.49 12.75 87.91 3.00 20.68 3.20 22.06 3.25 22.41 3.60 24.82			8	768 794 821 818 826 820 826 850	8.20 8.70 9.50 9.50 9.60 9.80 10.40	56.54 56.54 59.99 65.50 65.50 66.19 67.57 71.71			11	149 598 622 751 798 800 850 897 903	2.25 2.27 3.25 3.30 3.40	6.89 13.78 15.51 15.65 22.41 22.75 23.44
			437 458 471	4.00 27.57 4.30 29.65 4.80 33.10			0	449 568 580	3.00	19.99 20.69				962 970 1120	4.10	27.57 28.27 44.47
			493 498 505 515 517 531 548	5.00 34.47 5.20 35.86 5.30 36.54 5.40 37.23 5.70 39.30 5.75 39.65 6.20 42.75				631 663 684 749 755 777 781 781 850	4.30 4.80 5.70 6.60 6.60 6.70	26.89 29.65 33.10 39.30 41.36 45.51 45.51 46.20			12	732 800 850 883 946 1054 1247	2.25 2.90 4.10	6.89 11.38 13.79 15.51 19.99 28.27 44.47
			556 573 578 580 596 600	6.45 44.47 6.70 46.19 6.90 47.57 7.10 48.95 7.40 51.02 7.60 52.40				848 883 888 884 922	8.70 8.90 9.00 9.60	54.82 55.85 59.98 61.37 62.06 66.19			13	800 814 850 991 1182	0.65 .80 1.10 2.25 4.10	4.48 5.52 7.58 15.51 28.70
			615 620 619 621 616	7.90 54.47 8.00 55.15 8.00 55.15 8.10 55.85 8.10 55.85			9	915 937 537 648	10.00	9.31 19.99			14	1077 1140 1253 1520	2.90	15.51 19.99 28.27 47.58
			631 629 639	8.30 57.22 8.50 58.61 8.70 59.99				671 719 758	2.90 3.60 3.90	19.99 24.82 26.89			15	1029 1628	0.70	4.83 41.36
			635 645	8.78 60.54 9.10 62.75				759 804	4.75	28.27 32.75			16	1105	0.70	4.83
			675 685	10.00 68.94 10.00 68.94				902 927	7.00	44.13 48.26			18	1425	0.70	4.83

TABLE II.- MEASURED MODES, RESONANT FREQUENCIES; AND PRESSURES FOR ALL CYLINDERS - Continued

G-34 - 3 -	_]	f	P	g-14-13-		_	f		P	G-24::2	_ 1		f	1	р
Cylinder	m	n	cps or Hz	psi kn/m2	Cylinder	m	n	cps or Hz	psi	kN/m ²	Cylinder	m	n	cps or Hz	psi	kN/m ²
666	2	3	967	3.30 22.75	666	2	5	678 680	9.25	63.77	1001	1	8	821	6.00	41.36
			964 970 970	3.40 23.44 4.10 28.27 4.75 32.75				684 683 684	9.50	63.77 65.50 66.18 66.18			9	408 729	1.00	6.89 27.5
			978 978 973 982 982	5.40 37.23 5.80 39.99 6.30 43.44 6.35 43.78 6.50 44.81				693 696 714 723	10.00 10.20 10.90	68.94 70.32 75.15 76.53			10	474 610 816		6.89 13.78 27.5
			980 979 982 983	6.60 45.51 6.90 47.58 7.00 48.27 7.00 48.27			6	497 524 545	3.40 3.90	23.44 26.88 29.64			11	545 717 850		6.89 13.76 20.66
			990 991 986 992 991	8.00 55.15 8.10 55.85 8.20 56.54 8.85 61.02 8.90 61.37				544 591 607 609 624	4.40 5.25 5.70 6.00	30.33 36.19 39.30 41.36 42.74			12	617 784 951 1190	3.00	6.89 13.78 20.68 34.4
			990 997 997	9.60 66.19 9.60 66.19				630 646	6.45	44.47 47.57			13	657 1303	1.00 5.00	6.89 34.4
			995 993 1002 1054	9.80 67.57 9.90 68.26 10.00 68.95 11.11 76.60	1001	1	2	262 276 297 317	4.00	6.89 13.78 27.57 41.36			14	699 940 1260 1300	4.00	6.89 13.78 27.5 34.4
		4	666 670 682 677	3.00 20.69 3.40 23.44 3.50 24.13 4.30 29.65			3	262 280 297 313	3.00 4.00	13.78 20.68 27.57 34.47			15	798 1019 1568	1.00 2.00 5.00	6.89 13.70 34.4
			699 716 715 723	4.75 32.75 6.00 41.36 6.30 43.44 6.42 44.27			4	212 276 322	1.00	· ·			16	938 1169 1685		6.89 13.79 34.4
			716 725 739	6.70 46.20 6.95 47.92 8.10 55.85				401 430	5.00 6.00	34.47 41.36		2	3	316 336	4.00 5.00	13.70 27.5 34.4
			743 758 764 765 778 790	8.15 56.19 9.33 64.33 9.80 67.57 10.00 68.94 11.10 76.53 12.40 85.50			5	225 301 356 413 453 494	3.00 4.00 5.00	6.89 13.78 20.68 27.57 34.47 41.36			4	354 238 375 423	1.00	41.36 6.89 27.57 34.4
		5	537 545 556 554 579	3.30 22.75 3.65 25.17 3.90 26.89 4.00 27.57 4.00 27.57			6	283 365 443 507 517	1.00 2.00 3.00 4.00	<u> </u>			5	231 315 371 426 473 514	3.00 4.00 5.00	6.89 13.78 20.68 27.5 34.4 41.30
			570 581 579 585	4.30 29.65 4.80 33.20 4.80 33.09 5.00 34.47 5.95 41.02			7	309 421 719		6.89 13.78 41.36			6	287 520 576		6.8 27.5 34.4
			617 630 652 645 676	6.90 47.57 8.00 55.15 8.10 55.84 9.00 62.05			8	353 480 669 748	4.00	6.89 13.78 27.57 34.47			7	330 432 528 613	3.00	6.89 13.78 20.68 27.5

TABLE II.- MEASURED MODES, RESONANT FREQUENCIES, AND PRESSURES FOR ALL CYLINDERS - Continued

G-2 d-2			f		p	0-14-3-	_		f		р	0-34-3-	_		f	1	p
Cylinder	m	n	cps or Hz	psi	kn/m ²	Cylinder	m	n	cps or Hz	psi	kn/m²	Cylinder	m	n	cps or Hz	raq	kn/m²
1001	2	7	683	5.00	34.47	1502	2	7	550	5.00	34.47	1624	1	4	446 466		68.94
		8	373 608	1.00	6.89 20.68			10	506	5.00	13.78				481 494	12.00	75.84 82.73 89.63
			764		34.47	1624	1	2	424 419	1.00	6.89				514		96.52
1502	1	2	451 452		13.78 20.68				416 432	3.00	20.68			5	126 203	0	0 6.89
			459 461	4.00	27.57 34.47				420 433	4.00 5.00	27.57 34.47				266 281	2.30	13.78 15.85
		3	283	2.00	13.78				426 429	6.00	34.47 41.36				318 323	3.00	20.68 20.68
			296 314	4.00	20.68 27.57				448 430	7.00	47.57 48.26				349 370	4.00	24.13 27.57
		<u></u>	324	 	34.47				434 432	9.00	55.15 62.05				397 405	5.00	27.57 34.47
		4	240 2 96		13.78 27.57				432 433 444	11.00	68.94 75.84 75.84				429 463 472	6.00	34.47 41.36 41.36
	1	5	257 307		13.78 20.68	[[435 432	12.00	82.73 89.63				504 521	7.00	48.26 48.26
		<u></u>	378		34.47				421		96.52				505 521	8.00	55.15 62.05
		6	298 404		13.78 27.57			3	279 285	1.00	6.89 13.78				532 547		62.05
		<u> </u>	441		34.47				290 303	2.00	13.78 15.85				558 571	10.00	68.94
		7	342 411	3.00	13.80 20.68				309 317	3,50	20.68				586 594	11.00 12.00	82.73
			475		27•57				344 324	5.00	27.57 34.47				609 613	12.00 13.00	89.63
		8	392 472		13.78 20.68				334 346	6.00	41.36 41.36				627 665	13.00 14.00	96.52
		9	445 547	2.00	13.78				352 361 371	8.00	47.57 55.15 55.15			6	639	0	90.52
			616 687	4.00	27.57 34.47				369 380	9.00	62.05				156 301	1.00	
		10	· ·		13.78				382	10.00	68.94				321 368	2.00	13.78
			595 688		27.57				391 402	11.00	75.84 75.84				379 433 482	3.00 4.00	20.68
		11	551		13.78		ĺ		399 414	12.00	82.73 82.73				516	5.00	34.47
		13	653		13.78				418	13.00 13.00	89.63				525 568	7.00	41.36 48.26
		14	672	-	13.78					14.00 14.00					596 628	9.00	62.05
	2	3	647 659		20.68 27.57			4	170	0	0					9.00	68.94
		6	363	2.00	13.78				203 249	1.00 2.00 3.00	6.89 13.78				718	11.00	82.73
-			410 457 494	4.00	20.68 27.57 34.47				279 306 327	4.00 5.00	27.57		L			13.00 13.00	
	-	7	380	2.00					356 380	6.00	41.36 48.26			7	199 256	0.55	3.79 6.89
		_'	495	4.00	27.57				425	9.00					265	1.00	6.89

TABLE II.- MEASURED MODES, RESONANT FREQUENCIES, AND PRESSURES FOR ALL CYLINDERS - Concluded

1		Ţ	f	!	p				f		p
Cylinder	m	n	cps or Hz	psi	kN/m ²	Cylinder	m	n	cps or Hz	psi	kN/m ²
1624	1	7	316 364 382 441 479 497 524 551 558 606 617 634 652	1.46 2.00 2.2 3.00 3.5 4.00 4.4 5.00 5.00 6.00 6.3 7.00	10.06 13.78 15.16 20.68 24.13 27.57 30.53 34.47 34.36 41.36 43.43 48.26	1624	1	9	261 344 413 472 569 682 717 793 813 971 1017 1068	0.55 1.00 1.53 2.03 3.00 4.40 5.00 6.00 6.30 9.00 10.00	3.79 6.89 10.54 13.99 20.68 30.33 34.47 41.36 43.43 62.05 68.94 75.84
			740 775 814 820 881	9.00 10.00 11.00 11.00 13.00	62.05 68.94 75.84 75.84 89.63			10	237 294 389 458 530 529	0.25 .55 1.00 1.47 2.00 2.02	1.72 3.79 6.89 10.13 13.78 13.92
		8	231 289 358 412 419 437	0.55 1.00 1.47 1.94 2.00 2.20	3.79 6.89 10.13 13.37 13.78 15.16				569 632 631 696 772	2.38 3.00 3.00 3.50 4.44	16.40 20.68 20.68 24.13 30.61
			468 501 50 2 510 539 564	2.56 3.00 3.00 3.00 3.50 4.00	17.65 20.68 20.68 20.68 24.13 27.57			11	323 410 763 848 1293	0.55 1.00 3.50 4.40 11.00	3.79 6.89 24.13 30.33 75.84
			579 633 642 671 696	4.00 5.00 5.00 5.46 6.00	27.57 34.47 34.47 37.64 41.36			12	208 1004 1077 1223 1350	0 6.20 7.00 9.00 11.00	0 42.74 48.26 62.05 75.84
			709 751 807	6.10 7.00 8.00	42.05 48.26 55.15		2	3	672	9.00	62.05
			847 892	9.00	62.05 68.94			14	473 529	11.00	75.84 75.84
			943 1022	11.00	75.84 89.63			8	1183	14.00	96.52

TABLE III.- MEASURED FREQUENCIES AND MODES FOR MECHANICALLY APPLIED AXIAL LOADS FOR CYLINDERS 324 AND 666

(a) Cylinder 324

			Frequenc	y, f, cps (or H	z) for -	
m	n	ñ _x =0 n _x =0.3827x10 ⁻³ ñ _φ =0	$\bar{n}_{x}=0$ $n_{x}=0.1275\times10^{-3}$ $\bar{n}_{\phi}=0$	$\bar{n}_{x}=0$ $n_{x}=0.2551\times10^{-3}$ $\bar{n}_{\phi}=0$		n _x =0.0576×10 ⁻⁶ n _x =0.2551×10 ⁻³ n̄ _φ =0.4936×10 ⁻⁶
1	2345678	 254 177 167 196 247 323	329 250 172 162 191 243 320	399 252 175 165 194 246 321	262 216 231 276 346 415	407 281 266 303 364 527
2	5	303	291	294	337	400

m	n	f,* cps (or Hz) for $\bar{n}_{x}=n_{x}=\bar{n}_{\phi}=0$
1	a ₂ a ₃ a ₁ a ₅ a ₆ a ₇ a ₈ a ₉	337 228 168 160 190 244 317 399
2	5	284
3	ъ7	251

^{*}Aft ring disconnected from fixture.

a Node lines are on seam welds.

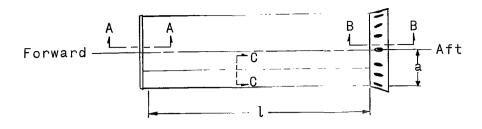
bNode lines are off seam welds.

TABLE III.- MEASURED FREQUENCIES AND MODES FOR MECHANICALLY APPLIED AXIAL LOADS FOR CYLINDERS 324 AND 666 - Concluded

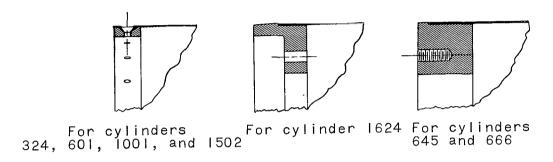
(b) Cylinder 666 with m = 1 and n = 3

$ar{\mathtt{n}}_{\mathbf{x}}$	π̄φ	Frequency, f, cps (or Hz) for $n_x = 0.8559 \times 10^{-3}$
0.248 × 10 ⁻⁶ .296 .307 .407	1.532 × 10 ⁻⁶ 1.829 1.898 2.515	470 485 492 526

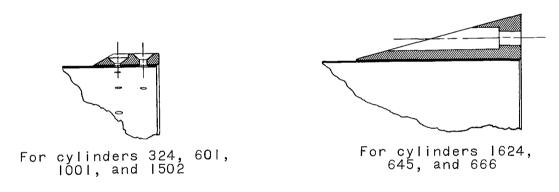
$ar{\mathtt{n}}_{\mathbf{x}}$	Frequency, f, cps (or Hz) for $\bar{n}_x = 0.299 \times 10^{-6}$ and $\bar{n}_\phi = 1.852 \times 10^{-6}$
0 -0.8559 × 10 ⁻³ -1.1984 -1.7120 -1.8833 -2.2599	481 480 468 466 452
Buck	le
-2.4312 -2.6365	438 344
Colla	pse



(a) Overall view of typical cylinder.



(b) Cross-sectional view of forward end-rings. Section AA of figure 1(a).

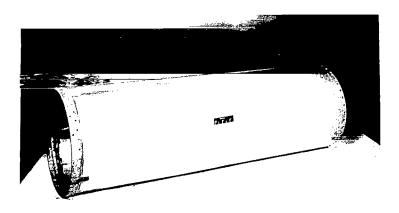


(c) Cross-sectional view of aft end-rings. Section BB of figure 1(a).



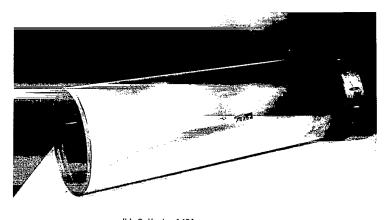
(d) Sketch of the 1/2-inch (1.27-cm) overlap used in the longitudinal seam joints. Section CC of figure 1(a).

Figure 1.- Details of cylinder.



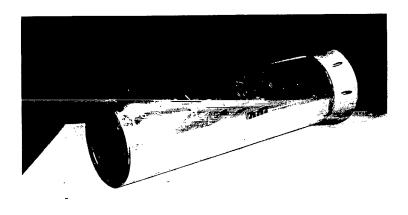
(a) Cylinder 1001.

L-62-8992



(b) Cylinder 1624.

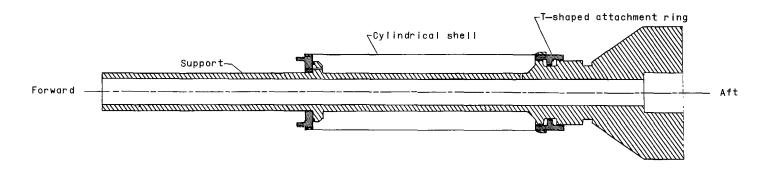
L-62-8989



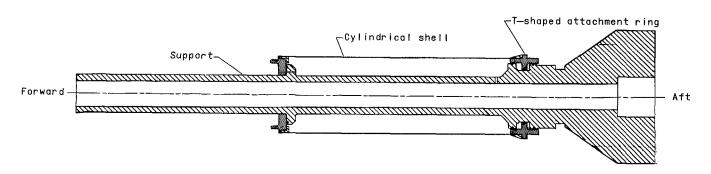
(c) Cylinder 645.

L-62-8988

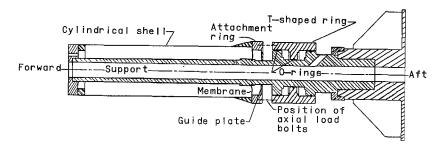
Figure 2.- Photographs of cylinders.



(a) Cylinder with 6-inch (15-cm) radius bolted radially to fixture.

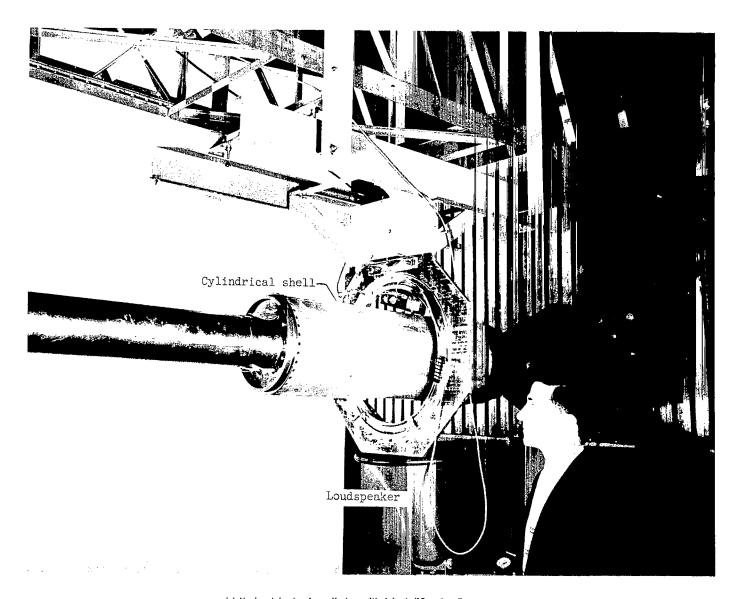


(b) Cylinder with 6-inch (15-cm) radius bolted longitudinally to fixture.



(c) Cylinder with 4-inch (10-cm) radius in place on fixture.

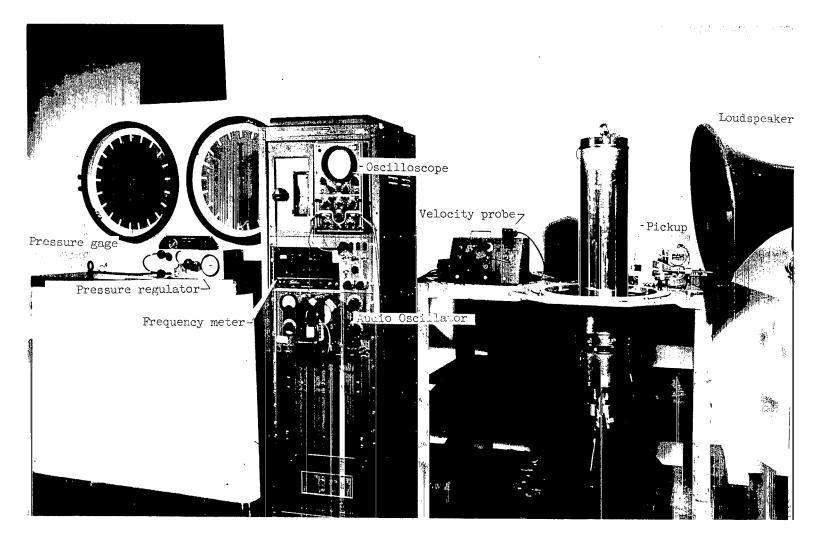
Figure 3.- Cross sections of cylinders.



(a) Horizontal setup for cylinders with 6-inch (15-cm) radius.

L-59-7777.1

Figure 4.- Photographs of test setups.



(b) Test equipment and vertical model setup.

L-62-8994.1

Figure 4.- Concluded.

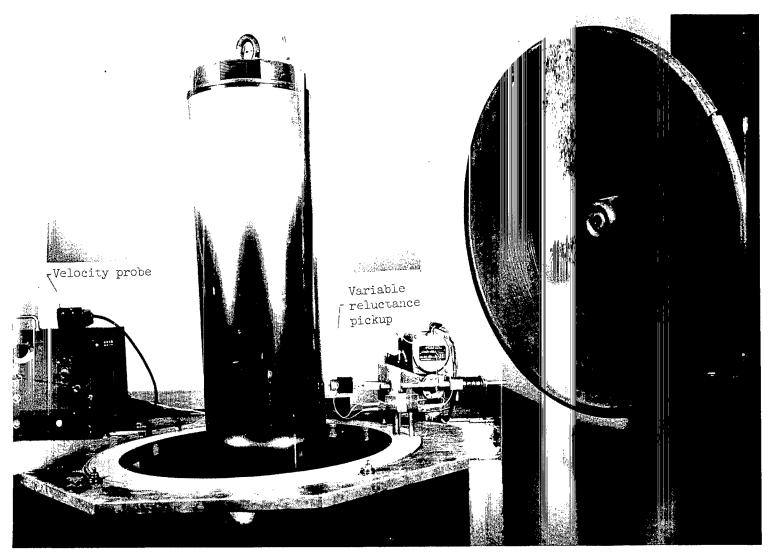


Figure 5.- Photograph of pickup.

L-62-8993.1

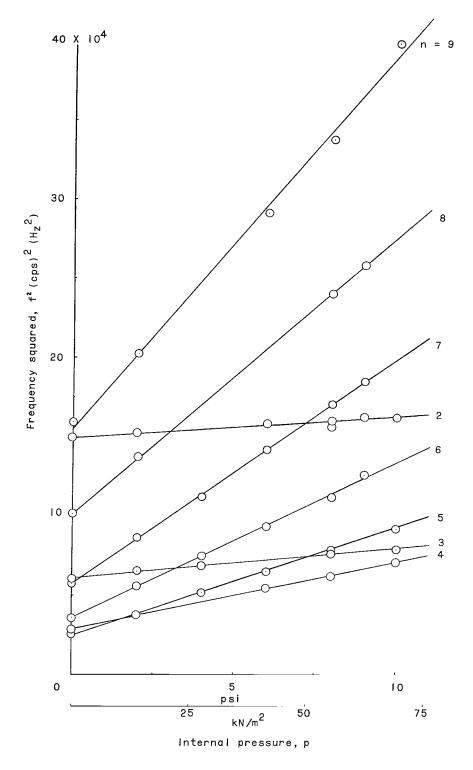


Figure 6.- Experimental variation of squared frequency with pressure for cylinder 324 with no external axial load.

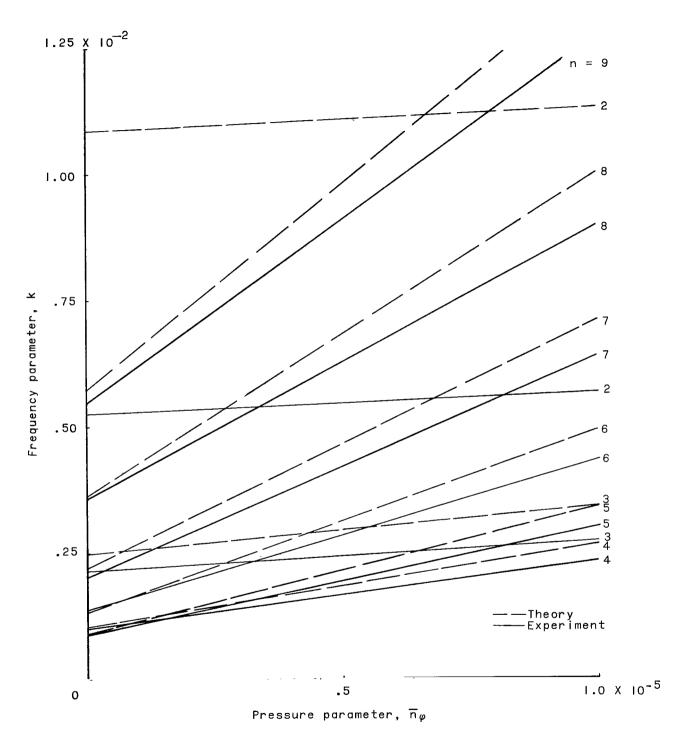


Figure 7.- Variation of frequency and pressure parameter with mode for cylinder 324. $n_{\chi}=0; m=1.$

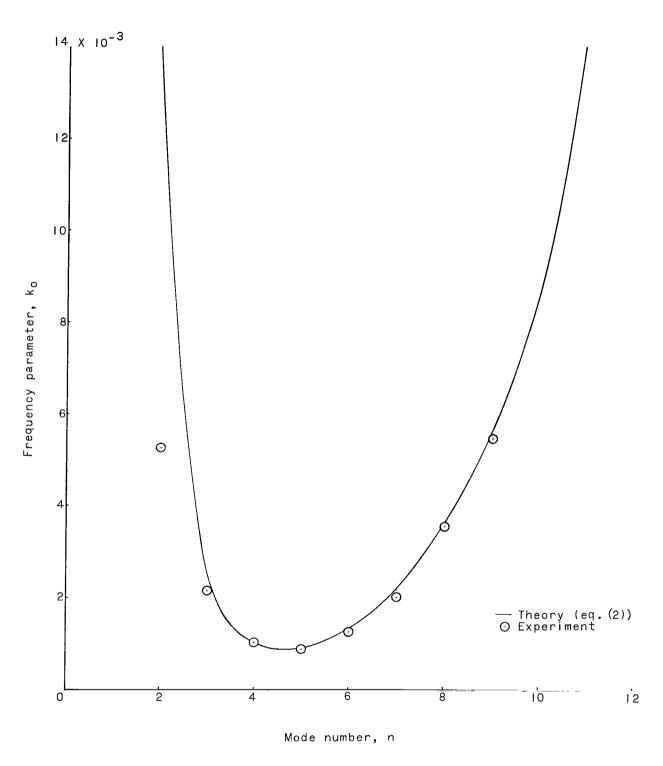


Figure 8.- Experimental and theoretical frequency parameter at zero load for various mode shapes for cylinder 324. m = 1.

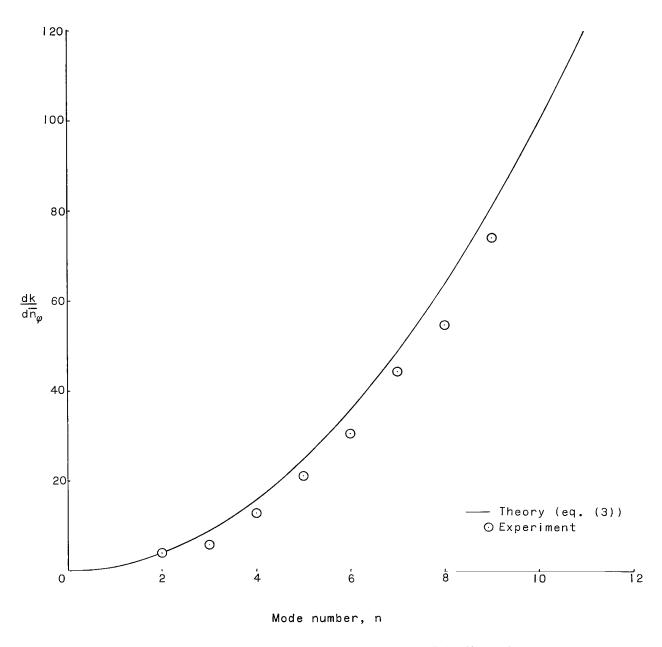


Figure 9.- Experimental and theoretical slope for various mode shapes for cylinder 324. m = 1.

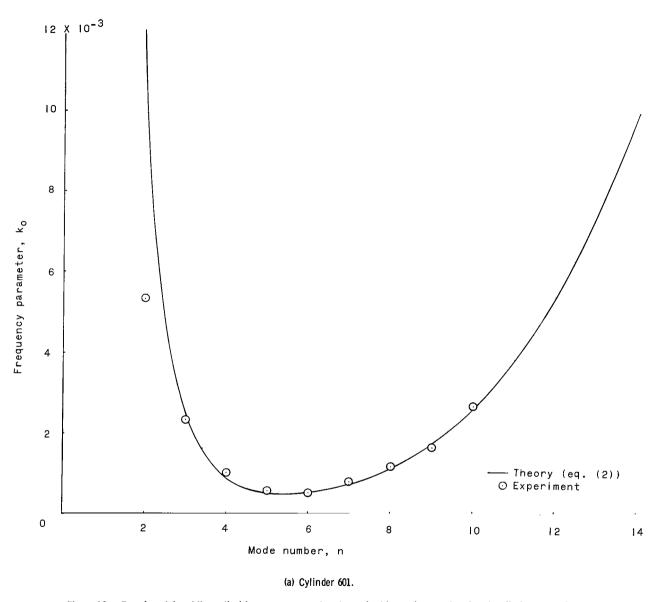


Figure 10.- Experimental and theoretical frequency parameter at zero load for various modes of each cylinder. m = 1.

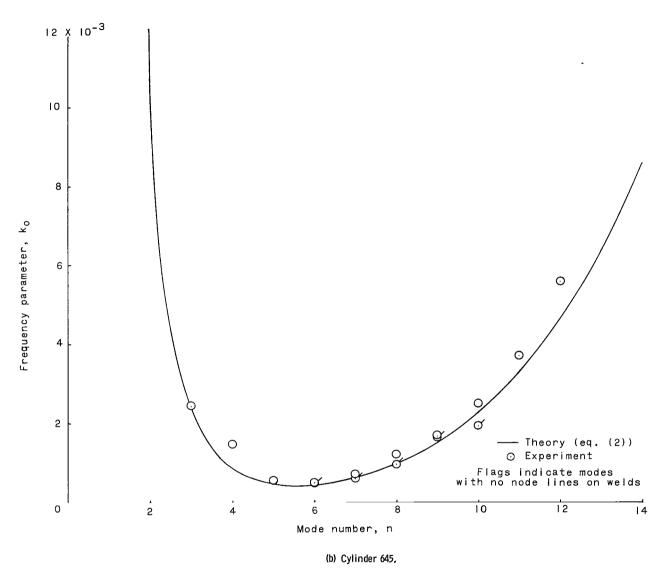


Figure 10.- Continued.

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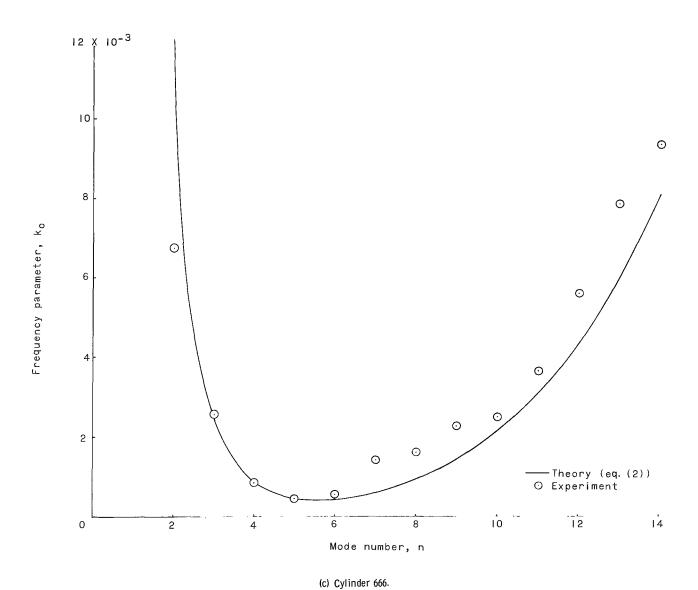
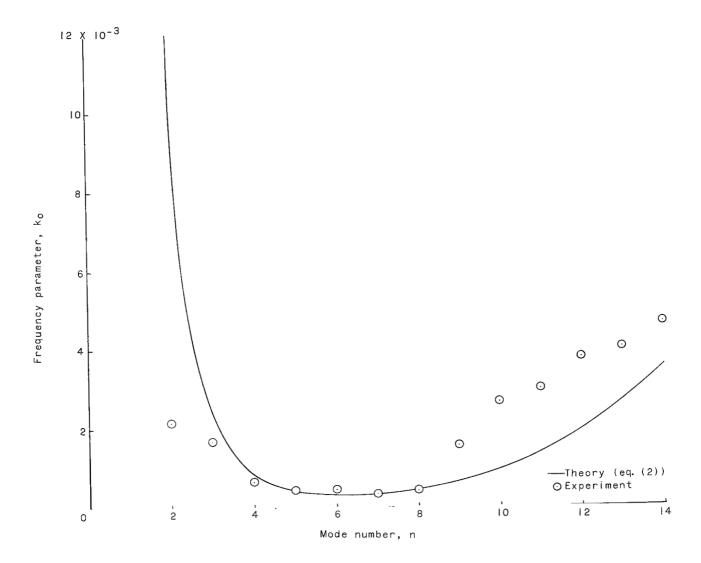


Figure 10.- Continued.



(d) Cylinder 1001.

Figure 10.- Continued.

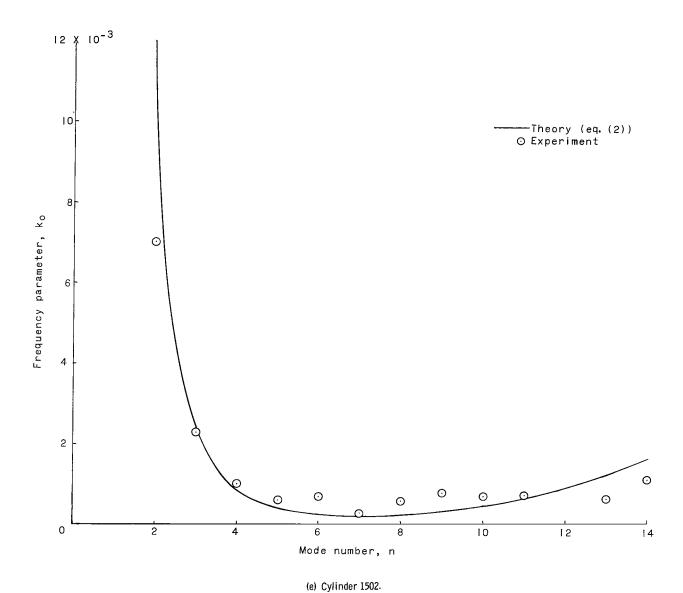


Figure 10.- Continued.

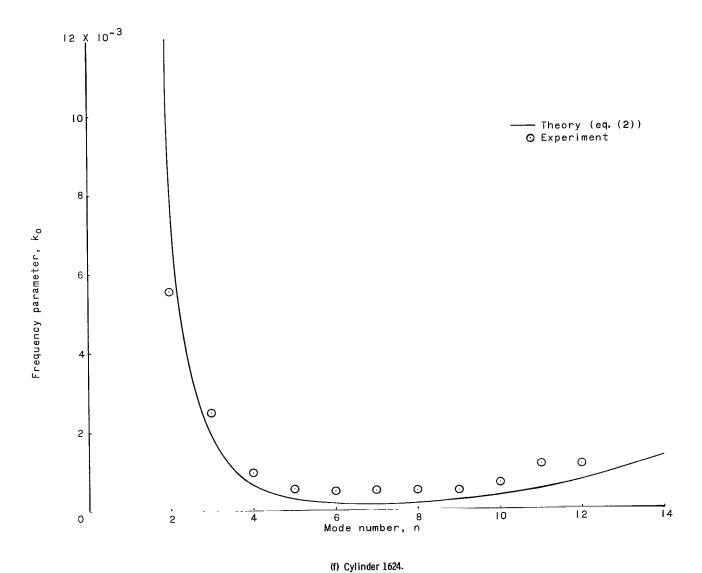


Figure 10.- Concluded.

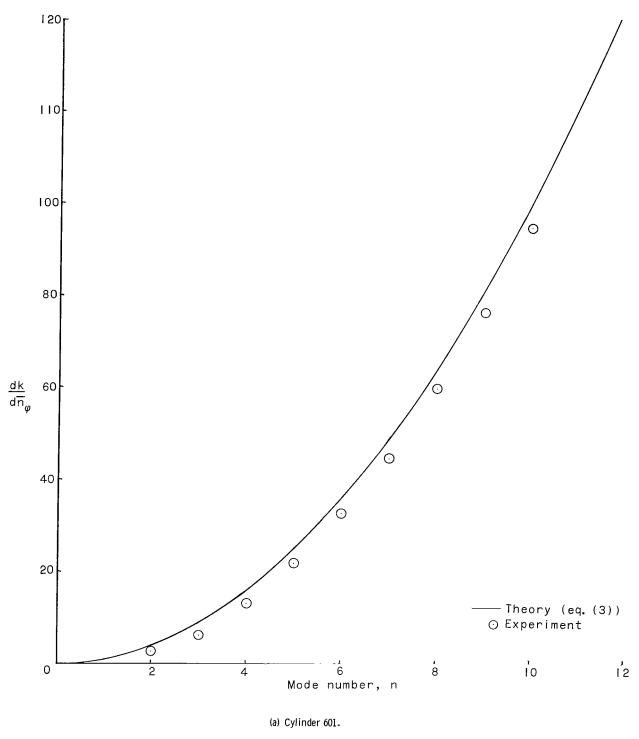
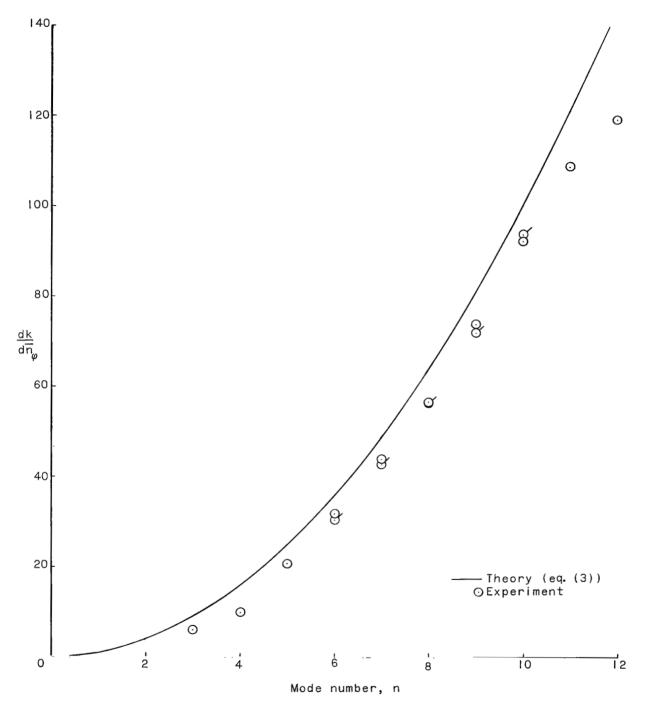


Figure 11.- Experimental and theoretical slope for various modes and cylinders. m = 1.



(b) Cylinder 645. Flags indicate modes with no node lines on welds.

Figure 11.- Continued.

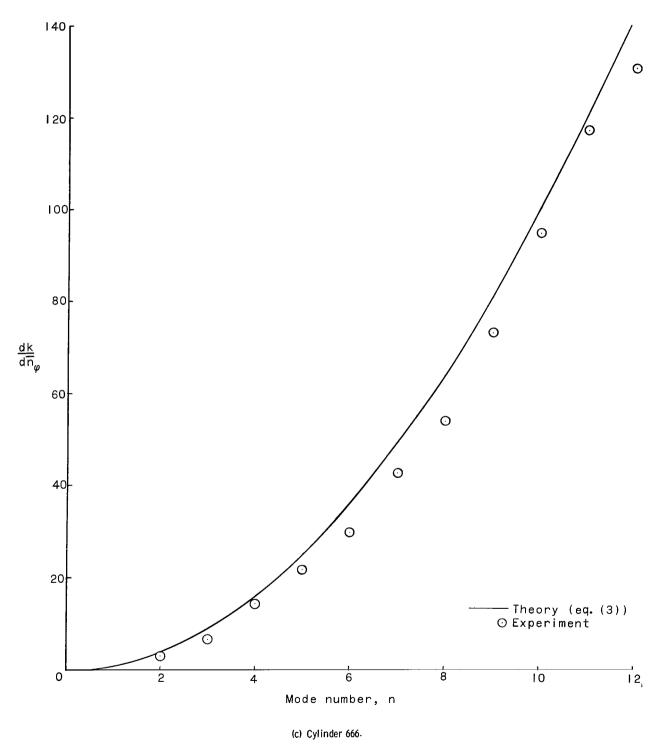


Figure 11.- Continued.

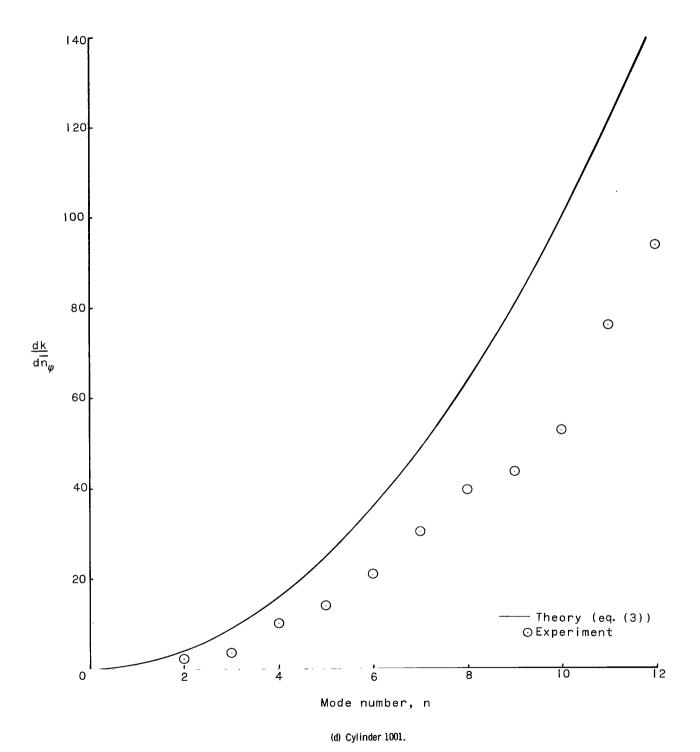


Figure 11.- Continued.

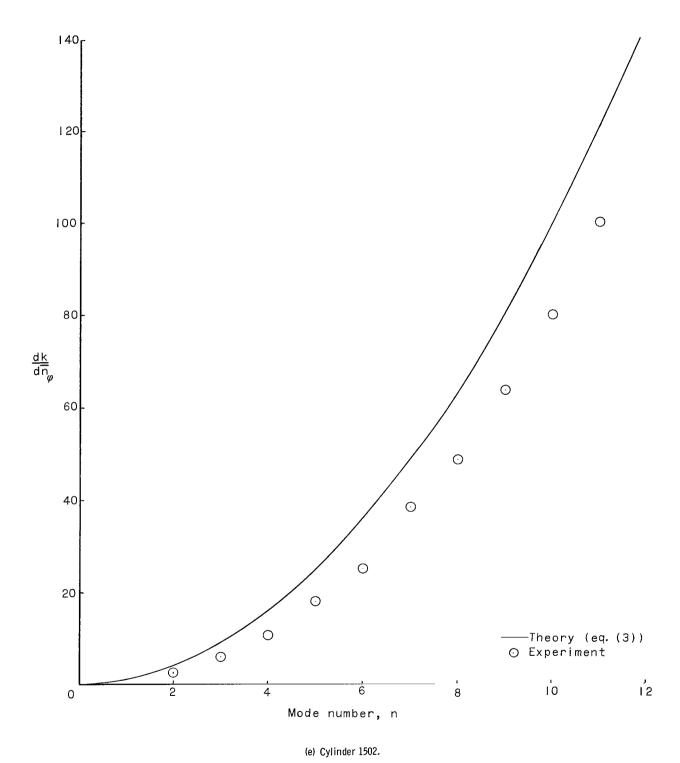


Figure 11.- Continued.

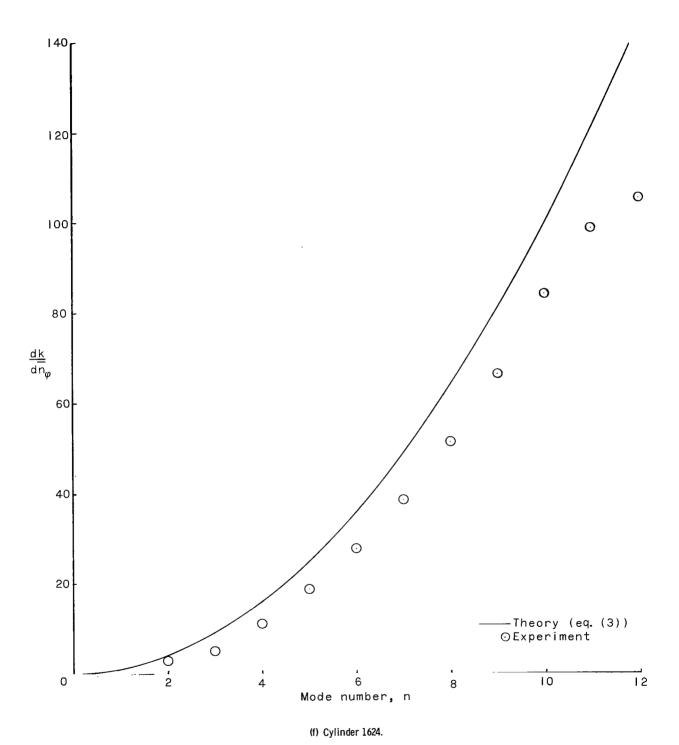


Figure 11.- Concluded.

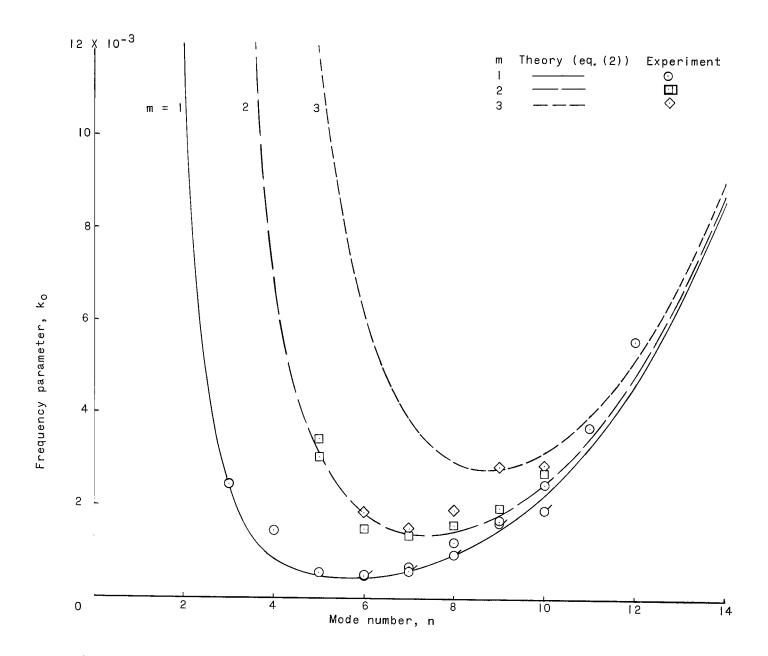


Figure 12.- Experimental and theoretical frequency parameter at zero load for various modes with different longitudinal wave numbers for cylinder 645.

Flags indicate modes with no node lines on welds.

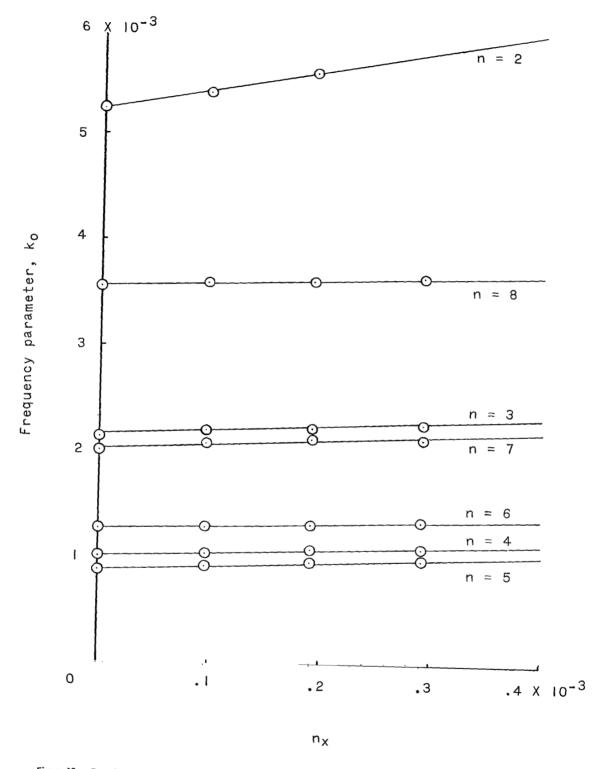


Figure 13.- Experimental variation of frequency parameter due to applied tension for cylinder 324. $\vec{n}_{\phi} = \vec{n}_{\chi} = 0$; m = 1.

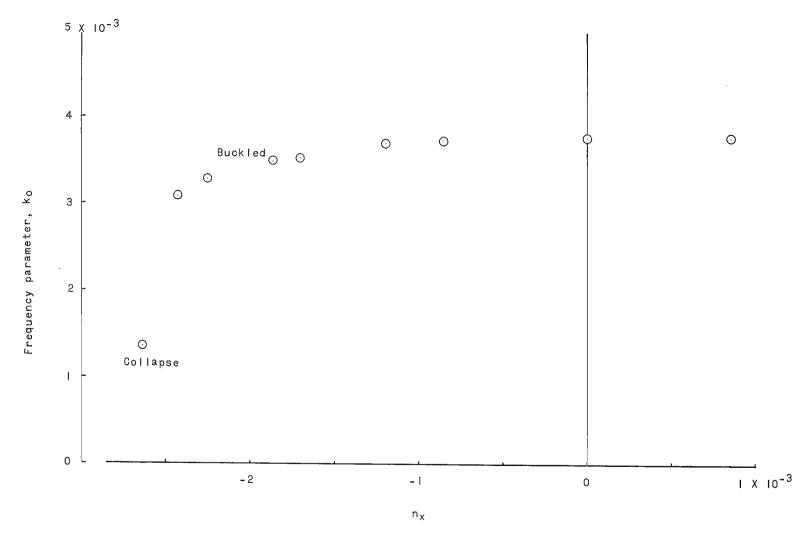


Figure 14.- Experimental frequency variation for mode m=1, n=3 through the buckling region for cylinder 666. $\overline{n}_{\varphi}=1.8\times10^{-6}$.

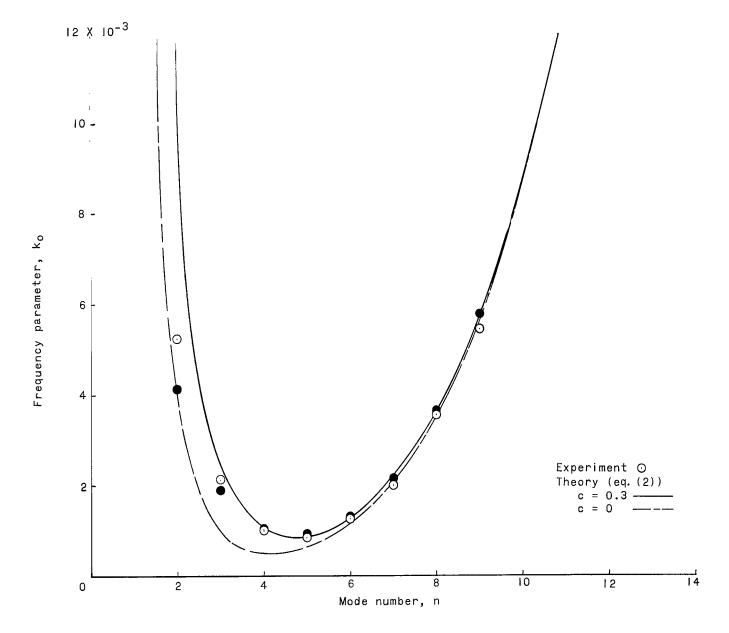


Figure 15.- Experimental and theoretical frequency parameter at zero load for various mode shapes including end fixity for cylinder 324. $n_{\chi} = 0$; m = 1. Dark symbols indicate that the aft ring was free.

3/18/8/5

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-NATIONAL AERONAUTICS AND SPACE ACT OF 1958

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